

### 193(4): Lagrangian Dynamics of Any Planar Orbit

The Lagrangian is:

$$L = \frac{1}{2} \mu (\dot{r}^2 + r^2 \dot{\theta}^2) - V(r) \quad - (1)$$

where  $\mu$  is the reduced mass:

$$\mu = \frac{m_1 m_2}{m_1 + m_2} \quad - (2)$$

and  $V(r)$  is the potential of attraction. If:

$$m_2 \gg m_1 \quad - (3)$$

the

$$\mu \approx m_1 \quad - (4)$$

to a good approximation. We denote by  $M$  the mass of the attracting object (e.g. the sun) and by  $m$  the mass of the attracted object (e.g. a planet). So:

$$\mu \approx m \quad - (5)$$

The Lagrangian is cyclically symmetric in  $\theta$  so the angular momentum  $L$  conjugate to the coordinate  $\theta$  is conserved:

$$\dot{L} = \frac{\partial L}{\partial \theta} = 0 = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} \quad - (6)$$

i.e.

$$L = \frac{\partial L}{\partial \dot{\theta}} = m r^2 \dot{\theta} \quad - (7)$$

= constant

The angular momentum  $L$  is the first integral of the motion, and is constant. It is the total angular momentum.

1) The Lagrange equation for  $r$  is:

$$\frac{\partial L}{\partial r} = \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} \quad - (8)$$

From eqs (1) and (8):

$$m(\ddot{r} - r\dot{\theta}^2) = -\frac{\partial V}{\partial r} = F(r) \quad - (9)$$

Here  $F(r)$  is the force of attraction

between  $m$  and  $M$ .

Make the change of variable:

$$u = 1/r \quad - (10)$$

then 
$$\frac{du}{d\theta} = -\frac{1}{r^2} \frac{dr}{d\theta} = -\frac{1}{r^2} \frac{dr}{dt} \frac{dt}{d\theta} \quad - (11)$$

From eq. (7):

$$\frac{d\theta}{dt} = \frac{L}{mr^2} \quad - (12)$$

so 
$$\frac{du}{d\theta} = -\frac{m}{L} \frac{dr}{dt} = -\frac{m}{L} \dot{r} \quad - (13)$$

and 
$$\frac{d^2 u}{d\theta^2} = \frac{d}{d\theta} \left( -\frac{m}{L} \dot{r} \right) = \frac{dt}{d\theta} \frac{d}{dt} \left( -\frac{m}{L} \dot{r} \right) = -\frac{m}{L} \ddot{r} \quad - (14)$$

Also: 
$$\frac{d^2 u}{d\theta^2} = -\frac{m^2}{L^2} r^2 \ddot{r} \quad - (15)$$

Therefore:

3)

$$\ddot{r} = -\frac{L^2}{m^2} u^2 \frac{d^2 u}{d\theta^2} - (16)$$

$$r \ddot{\theta} = \frac{L^2}{m^2} u^3 - (17)$$

So

$$\frac{d^2 u}{d\theta^2} + u = -\frac{m}{L^2} \frac{1}{u^2} F(u) - (18)$$

and the general force law for any planar orbit is:

$$F(r) = -\frac{L^2}{mr^2} \left[ \frac{d^2}{d\theta^2} \left( \frac{1}{r} \right) + \frac{1}{r} \right] - (19)$$

This is true for any potential of attraction defined by:

$$F(r) = -\frac{\partial V(r)}{\partial r} - (20)$$

### The Static Conic Sections

These are all defined by:

$$\frac{1}{r} = \frac{1}{d} (1 + e \cos \theta) - (21)$$

where  $d$  is the half right latitude and  $e$  is the eccentricity. The ellipse is defined by  $0 < e < 1$ , the hyperbola by  $e > 1$ , the parabola by  $e = 1$ , and the circle by  $e = 0$ .

4) For all the conic sections the force law is:

$$F(r) = -\frac{L^2}{m d r^3} \quad - (22)$$

and eccentricity  $e$  does not appear in the force law.

For the circle for example:

$$d = r \quad - (23)$$

so for a circular orbit:

$$F(r) = -\frac{L^2}{m r^3} \quad - (24)$$

which is an inverse cube force law. From eq:

$$(22): \quad F(u) = -\left(\frac{L^2}{m d}\right) u^2 \quad - (25)$$

so for eqs. (18) and (25):

$$\frac{d^2 u}{d\theta^2} + u = \frac{1}{d} \quad - (26)$$

for all conic sections, i.e.

$$\boxed{\frac{d^2}{d\theta^2} \left( \frac{1}{r} \right) + \frac{1}{r} = \frac{1}{d}} \quad - (27)$$

The conic sections we sketch as follows:

5)

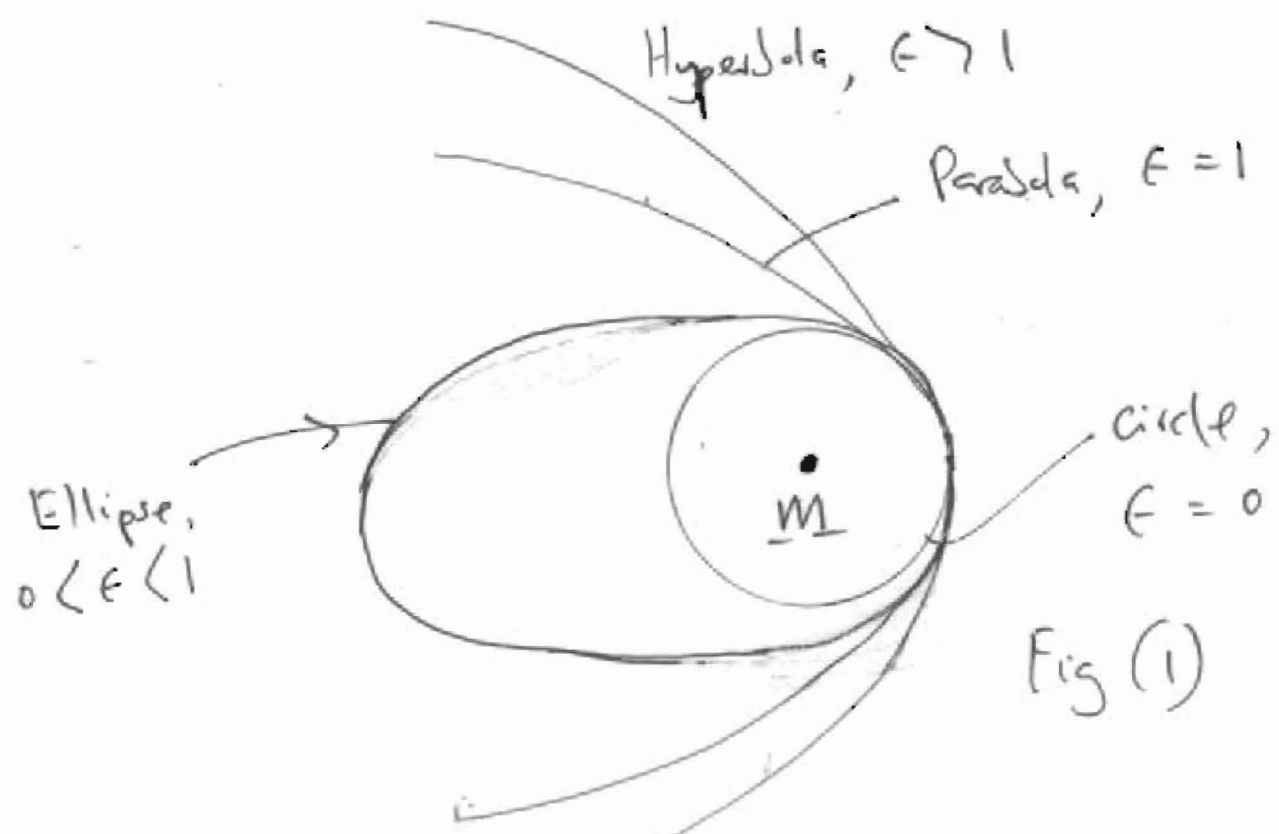


Fig (1)

So the deflection of an object by the object  $M$  can be calculated from eq. (27) or eq. (26). The deflection does not depend on  $e$ .

The Precessing Conic Sections

In this case:

$$\frac{1}{r} = \frac{1}{d} (1 + e \cos(x\theta)) \quad (28)$$

and

$$F(r) = -\frac{L^2}{mr^2} \left( \frac{x^2}{d} + \frac{1}{r} (1 - x^2) \right) \quad (29)$$

so

$$F(u) = -\frac{L^2 x^2}{md} u^2 - \frac{L^2 (1 - x^2)}{m} u^3 \quad (30)$$

and

$$\boxed{\frac{d^2 u}{d\theta^2} + u = \frac{x^2}{d} + (1 - x^2) u} \quad (31)$$