

Note 320(b): Complete Calculation of the Gravitomagnetic Field.

The calculation begins with the general coordinate transform:

$$\Gamma^{\mu'\nu'} = \frac{\partial x^{\mu'}}{\partial x^{\mu}} \frac{\partial x^{\nu'}}{\partial x^{\nu}} \Gamma^{\mu\nu} \quad (1)$$

of the gravitomagnetic field tensor. Here $\Gamma^{\mu\nu}$ is the rest frame, a frame in which the particle is at rest, and $\Gamma^{\mu'\nu'}$ the frame in which the particle moves with a velocity \underline{v} . This is the observer frame. For example, if the observer is situated on the sun, the earth orbits the sun with velocity \underline{v} .

Eq. (1) means that:

$$\begin{aligned} \underline{g}' &= \gamma (\underline{g} + \underline{v} \times \underline{\Omega}) - \frac{\gamma^2}{1+\gamma} \frac{\underline{v}}{c} \left(\frac{\underline{v}}{c} \cdot \underline{g} \right) \\ \underline{\Omega}' &= \gamma \left(\underline{\Omega} - \frac{1}{c^2} \underline{v} \times \underline{g} \right) - \frac{\gamma^2}{1+\gamma} \frac{\underline{v}}{c} \left(\frac{\underline{v}}{c} \cdot \underline{\Omega} \right) \end{aligned} \quad (2)$$

where

$$\gamma = \left(1 - \frac{v^2}{c^2} \right)^{-1/2} \quad (4)$$

A Lorentz transform has been used because of the pseudo Lorentz transform property of the gravitomagnetic field equations.

2) In the limit: $\gamma \ll 1$ - (5)

$$\underline{g}' = \underline{g} + \underline{v} \times \underline{\Omega} \quad - (6)$$

$$\underline{\Omega}' = \underline{\Omega} - \frac{1}{c^2} \underline{v} \times \underline{g} \quad - (7)$$

The analogous equations in electrodynamics are:

$$\underline{E}' = \underline{E} + \underline{v} \times \underline{B} \quad - (8)$$

$$\underline{B}' = \underline{B} - \frac{1}{c^2} \underline{v} \times \underline{E} \quad - (9)$$

In cylindrical polar coordinates the linear velocity is:

$$\underline{v} = \dot{r} \underline{e}_r + \omega r \underline{e}_\theta \quad - (10)$$

and the linear acceleration is:

$$\underline{a} = (\ddot{r} - r\dot{\theta}^2) \underline{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{e}_\theta \quad - (11)$$

In Cartesian coordinates:

$$\underline{v} = \dot{r} \underline{e}_r \quad - (12)$$

and

$$\underline{a} = \ddot{r} \underline{e}_r \quad - (13)$$

In cylindrical polar coordinates the frame is rotating, and generates the extra terms.

3) In vector notation:

$$\underline{v} = \dot{r} \underline{e}_r + \underline{\omega} \times \underline{r} \quad - (14)$$

and:

$$\underline{a} = \ddot{r} \underline{e}_r - \underline{\omega} \times (\underline{\omega} \times \underline{r}) + \frac{d\underline{\omega}}{dt} \times \underline{r} + 2\underline{\omega} \times \frac{dr}{dt} \underline{e}_r \quad - (15)$$

where

$$\underline{e}_r = \underline{e}_\theta \times \underline{k} \quad - (16)$$

$$\underline{e}_\theta = \underline{k} \times \underline{e}_r \quad - (17)$$

$$\underline{k} = \underline{e}_r \times \underline{e}_\theta \quad - (18)$$

In gravitational theory,

$$\underline{a} = \underline{g} \quad - (19)$$

$$\text{so } \underline{g} = \ddot{r} \underline{e}_r - r \dot{\theta}^2 \underline{e}_r + (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) \underline{e}_\theta \quad - (20)$$

In this notation, $\underline{e}_r(\theta)$ is:

$$\underline{g} = \ddot{r} \underline{e}_r + \underline{v} \times \underline{\Omega} \quad - (21)$$

where

$$\underline{v} = \dot{r} \underline{e}_r + \underline{\omega} \times \underline{r} \quad - (22)$$

therefore in general:

$$\underline{v} \times \underline{\Omega} = -r \dot{\theta}^2 \underline{e}_r + (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) \underline{e}_\theta \quad - (23)$$

4) Planar Orbital Theory

In this case:

$$(r\ddot{\theta} + 2\dot{r}\dot{\theta})\underline{e}_{\theta} = \underline{0} \quad - (24)$$

for any planar orbit. The result (24) follows from Lagrangian dynamics as in previous work. So for planar orbits:

$$\underline{v} = \dot{r}\underline{e}_r + \underline{\omega} \times \underline{r} \quad - (25)$$

and

$$\underline{a} = \ddot{r}\underline{e}_r - \underline{\omega} \times (\underline{\omega} \times \underline{r}) \quad - (26)$$

From eqs. (21) and (26):

$$\underline{v} \times \underline{\Omega} = -\underline{\omega} \times (\underline{\omega} \times \underline{r}) \quad - (27)$$

The solution of eq. (27) is obtained from:

$$(\dot{r}\underline{e}_r + \underline{\omega} \times \underline{r}) \times \underline{\Omega} = -\underline{\omega} \times (\underline{\omega} \times \underline{r}) \quad - (28)$$

The solution of eq. (28) is:

$$\begin{aligned} (\underline{\omega} \times \underline{r}) \times \underline{\Omega} &= -\underline{\Omega} \times (\underline{\omega} \times \underline{r}) \quad - (29) \\ &= -\underline{\omega} \times (\underline{\omega} \times \underline{r}) \end{aligned}$$

i.e.

$$\underline{\Omega} = \underline{\omega} \quad - (30)$$

and

$$\underline{v} = \underline{\omega} \times \underline{r} \quad - (31)$$

Therefore in planar orbital theory the

5) gravitomagnetic field is:

$$\underline{\Omega} = \underline{\omega} = \frac{d\theta}{dt} \underline{k} \quad - (32)$$

which is the angular velocity of the rotating frame.
 For example $\underline{\Omega}$ is the angular velocity of the earth rotating about its axis. The velocity \underline{v} of the Lorentz transform is the orbital linear velocity generated by the angular velocity of the rotating frame.

This is a simple but profound result, indicating that planetary orbital dynamics is the result of the EFE2 gravitomagnetic field then, a generally covariant unified field theory. This is a natural way of describing orbital dynamics with general relativity. In general if dynamics the Coriolis force is non-zero, so the gravitomagnetic field is in general:

$$\underline{\Omega} = \underline{\Omega}_1 \quad - (33)$$

Now assume that:

$$\underline{v} \times \underline{\Omega}_1 = (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{e}_\theta \quad - (34)$$

The solution is

$$\underline{v} = \frac{dr}{dt} \underline{e}_r \quad - (35)$$

and

$$\frac{dr}{dt} \underline{e}_r \times \underline{\Omega}_1 = (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{e}_\theta \quad - (36)$$

6) Using:

$$\underline{e}_\theta = \underline{k} \times \underline{e}_r \quad - (37)$$

it follows that:

$$\underline{\Omega}_1 = - \frac{1}{\dot{r}} (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{k} \quad - (38)$$

$$= - \left(\frac{r}{\dot{r}} \ddot{\theta} + 2\dot{\theta} \right) \underline{k} \quad - (39)$$

and:

$$\underline{a} = \ddot{r} \underline{e}_r + (\underline{\omega} \times \underline{r}) \times \underline{\Omega}_0 + \frac{dr}{dt} \underline{e}_r \times \underline{\Omega}_1$$

where:

$$\underline{\Omega}_0 = \omega \underline{k} \quad - (40)$$

and

$$\underline{\Omega}_1 = - \left(\frac{r}{\dot{r}} \ddot{\theta} + 2\dot{\theta} \right) \underline{k} \quad - (41)$$

Therefore in general, electrodynamics are described by gravitomagnetic Lorentz force equations.
