

### 401(3): Simple ECE2 Theory of Orbital Precession

Consider the fundamental definition of force in ECE2 gravitational theory:

$$\underline{F} = -\underline{\nabla} \underline{\Phi} + \underline{\omega} \underline{\Phi} \quad - (1)$$

where  $\underline{\Phi}$  is the gravitational potential and  $\underline{\omega}$  the spin correction. The vacuum force is:

$$\underline{F}(\text{vac}) = \underline{\omega} \underline{\Phi} \quad - (2)$$

where  $\underline{\Phi}$  is the gravitational potential in the absence of the vacuum:

$$\underline{\Phi} = -\frac{mG}{r} \quad - (3)$$

In Cartesian coordinates eq. (1) is:

$$\ddot{x} = -\frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{x}{x^2+y^2} + \omega_x \right) \quad - (4)$$

$$\ddot{y} = -\frac{mG}{(x^2+y^2)^{3/2}} \left( \frac{y}{x^2+y^2} + \omega_y \right) \quad - (5)$$

From the formal Taylor expansion of the previous note:

$$\omega_x \underline{\Phi} = \langle F_x(\text{vac}) \rangle^{(2)} + \langle F_x(\text{vac}) \rangle^{(4)} + \dots \quad - (6)$$

$$\omega_y \underline{\Phi} = \langle F_y(\text{vac}) \rangle^{(2)} + \langle F_y(\text{vac}) \rangle^{(4)} + \dots \quad - (7)$$

in which  $\underline{\Phi}$  is given by eq. (3).

In ECE2 theory the gravitational force (1) is equivalent to the Lagrangian:

2)

$$L = -\frac{mc^2}{\gamma} + \frac{mMG}{r} \quad - (8)$$

and the Hamiltonian:

$$H = \gamma mc^2 - \frac{mMG}{r} \quad - (9)$$

and the field equations of <sup>E(2)</sup> gravitation, written in a space with finite torsion and curvature  
Forward precession is produced by the Euler

Lagrange equations:

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} \quad - (10)$$

$$\frac{\partial L}{\partial y} = \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} \quad - (11)$$

and the backward precession is produced by the relativistic  
Newton equation:

$$\gamma^3 m \ddot{\underline{r}} = -mMG \frac{\underline{r}}{r^3} \quad - (12)$$

which is obtained from the Euler Lagrange equation:

$$\frac{\partial L}{\partial \underline{r}} = \frac{d}{dt} \frac{\partial L}{\partial \dot{\underline{r}}} \quad - (13)$$

and Lagrangian (8).

As in Note 401 (1), forward precession is given  
for eqs. (10) and (11) by:

$$\ddot{X} = \frac{mG}{\gamma (x^2 + y^2)^{3/2}} \left( \frac{\dot{X}\dot{Y} + X\dot{X}^2}{c^2} - X \right) \quad - (14)$$

$$\ddot{Y} = \frac{mG}{\gamma (x^2 + y^2)^{3/2}} \left( \frac{\dot{Y}\dot{X} + Y\dot{Y}^2}{c^2} - Y \right) \quad - (15)$$

Spic corrections for forward precession are found by a comparison of eqns. (4) and (5) and (14) and (15). In these equations:

$$\frac{1}{\gamma} = \left( 1 - \frac{\dot{X}^2 + \dot{Y}^2}{c^2} \right)^{1/2} \quad - (16)$$

For retrograde precession:

$$\ddot{X} = -\frac{mG}{\gamma^3} \frac{X}{(x^2 + y^2)^{3/2}} \quad - (17)$$

$$\ddot{Y} = -\frac{mG}{\gamma^3} \frac{Y}{(x^2 + y^2)^{3/2}} \quad - (18)$$

Spic corrections for retrograde precession are found by a comparison of eqns. (4) and (5) and (17) and (18).

These spic correction components can now be used in eqns. (6) and (7). In previous work it was shown that

$$\langle F_x(\text{vac}) \rangle^{(2)} = \langle F_y(\text{vac}) \rangle^{(2)} = 0 \quad - (19)$$

Assuming that the tensorial Taylor series converge rapidly, then:

$$\omega_x \Phi = \langle F_x(\text{vac}) \rangle^{(4)} - (20)$$

$$\omega_y \Phi = \langle F_y(\text{vac}) \rangle^{(4)} - (21)$$

in which  $\Phi$ ,  $\omega_x$  and  $\omega_y$  are known. It is therefore possible to find  $\langle (\underline{\delta r} \cdot \underline{\delta r})^2 \rangle$  became:

$$\langle F_x(\text{vac}) \rangle^{(4)} = \frac{1}{216} \langle (\underline{\delta r} \cdot \underline{\delta r})^2 \rangle \left( \frac{\partial^4 F_x}{\partial x^4} + \frac{\partial^4 F_x}{\partial y^4} + 6 \frac{\partial^4 F_x}{\partial x^2 \partial y^2} \right) - (22)$$

$$\text{and } \langle F_y(\text{vac}) \rangle^{(4)} = \frac{1}{216} \langle (\underline{\delta r} \cdot \underline{\delta r})^2 \rangle \left( \frac{\partial^4 F_y}{\partial x^4} + \frac{\partial^4 F_y}{\partial y^4} + 6 \frac{\partial^4 F_y}{\partial x^2 \partial y^2} \right) - (23)$$

it can also be in the  $xy$  plane.

It would be very interesting to graph the vacuum fluctuation  $\langle (\underline{\delta r} \cdot \underline{\delta r})^2 \rangle$  far forward and retrograde precession, showing that the origin of ECE2 relativity is vacuum fluctuation on the classical level. The overall aim of the calculation is to reproduce observed orbital precession with precision.